

A BERWALD Γ_0 – LINEAR CONNECTION $B\Gamma_0$ IN THE RIEMANN-LAGRANGE GEOMETRY OF 1-JET SPACES

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Abstract: This paper introduces the notions of a nonlinear connection Γ and of a Γ -linear connection $\nabla\Gamma$ on the 1-jet space $J^1(T, M)$. A particular nonlinear connection Γ_0 and a Berwald Γ_0 -linear connection $B\Gamma_0$ are produced by a pair of semi-Riemannian metrics. The adapted components of the torsion and curvature d-tensors of our Berwald connection are described.

Key words: 1-jet spaces, nonlinear connections, Γ -linear connections, a Berwald connection.

1. INTRODUCTION

According to Olver’s opinion expressed in the monograph [4] and in some private discussions, we emphasize that the 1-jet bundle represents the most convenient *space of configurations* for the study of quantum and classical field theories.

For that reason, many researchers studied the differential geometry of the 1-jet spaces, in the sense of d-connections, d-torsions and d-curvatures.

The geometrical approach from this paper follows the direction of development of the differential geometry of the 1-jet space $J^1(T, M)$ initiated by Asanov [1] and uses the geometrical methods from the theory of Lagrange spaces developed by Miron and Anastasiei [2].

It is important to note that this geometrical approach allows a clear exposition of the multi-time physical-mathematical concepts studied by Neagu [3] and offers many original ideas for the geometric dynamics of a PDEs systems, developed by Udriște [5].

2. NONLINEAR CONNECTIONS

Let us consider the jet bundle of first order $J^1(T, M) \rightarrow T \times M$, whose local coordinates $(t^\alpha, x^i, x^i_\alpha)$, where $\alpha = \overline{1, p}$, $i = \overline{1, n}$, transform after the rules:

$$\begin{aligned} \tilde{t}^\alpha &= \tilde{t}^\alpha(t^\beta), \quad \tilde{x}^i = \tilde{x}^i(x^j), \\ \tilde{x}^i_\alpha &= \frac{\partial \tilde{x}^i}{\partial x^j} \frac{\partial t^\beta}{\partial \tilde{t}^\alpha} x^j_\beta \end{aligned} \quad (1)$$

By definition, a pair of local functions $\Gamma = (M^{(i)}_{(\alpha)\beta}, N^{(i)}_{(\alpha)j})$, which transform after the rules:

$$\begin{aligned} \tilde{M}^{(j)}_{(\beta)\mu} \frac{\partial \tilde{t}^\mu}{\partial t^\alpha} &= M^{(k)}_{(\gamma)\alpha} \frac{\partial \tilde{x}^j}{\partial x^k} \frac{\partial t^\gamma}{\partial \tilde{t}^\beta} - \frac{\partial \tilde{x}^j_\beta}{\partial t^\alpha} \\ \tilde{N}^{(j)}_{(\beta)k} \frac{\partial \tilde{x}^k}{\partial x^i} &= N^{(k)}_{(\gamma)i} \frac{\partial \tilde{x}^j}{\partial x^k} \frac{\partial t^\gamma}{\partial \tilde{t}^\beta} - \frac{\partial \tilde{x}^j_\beta}{\partial x^i} \end{aligned} \quad (2)$$

is called a *nonlinear connection* on the 1-jet bundle $E = J^1(T, M)$.

A nonlinear connection Γ on the 1-jet space E produces an *adapted basis* of vector fields:

$$\left\{ \frac{\delta}{\delta t^\alpha}, \frac{\delta}{\delta x^i}, \frac{\partial}{\partial x_\alpha^i} \right\} \subset X(E) \quad (3)$$

where

$$\begin{aligned} \frac{\delta}{\delta t^\alpha} &= \frac{\partial}{\partial t^\alpha} - M_{(\beta)\alpha}^{(j)} \frac{\partial}{\partial x_\beta^j} \\ \frac{\delta}{\delta x^i} &= \frac{\partial}{\partial x^i} - N_{(\beta)i}^{(j)} \frac{\partial}{\partial x_\beta^j} \end{aligned} \quad (4)$$

This adapted basis is extremely convenient in the study of the differential geometry of 1-jet spaces because the transformation rules of its elements have a simple tensorial form:

$$\begin{aligned} \frac{\delta}{\delta t^\alpha} &= \frac{\partial \tilde{t}^\beta}{\partial t^\alpha} \frac{\delta}{\delta \tilde{t}^\beta} \\ \frac{\delta}{\delta x^i} &= \frac{\partial \tilde{x}^j}{\partial x^i} \frac{\delta}{\delta \tilde{x}^j} \\ \frac{\partial}{\partial x_\alpha^i} &= \frac{\partial \tilde{x}^j}{\partial x^i} \frac{\partial t^\alpha}{\partial \tilde{t}^\beta} \frac{\partial}{\partial \tilde{x}_\beta^j} \end{aligned} \quad (5)$$

In this context, the Lie algebra $X(E)$ of the vector fields on the 1-jet bundle E decomposes as the following direct sum:

$$X(E) = X(H_T) \oplus X(H_M) \oplus X(V)$$

where

$$\begin{aligned} X(H_T) &= \text{Span} \left\{ \frac{\delta}{\delta t^\alpha} \right\} \\ X(H_M) &= \text{Span} \left\{ \frac{\delta}{\delta x^i} \right\} \\ X(V) &= \text{Span} \left\{ \frac{\partial}{\partial x_\alpha^i} \right\} \end{aligned} \quad (6)$$

As a consequence, any vector field $X \in X(E)$ on the 1-jet space E can be unique written in the form:

$$X = h_T X + h_M X + v X \quad (7)$$

where h_T, h_M and v are the *canonical projections* of the above decomposition of the Lie algebra $X(E)$.

3. Γ -LINEAR CONNECTIONS

Let Γ be a nonlinear connection on the 1-jet space $E = J^1(T, M)$ and let h_T, h_M and v be the canonical projections of the decomposition of the Lie algebra of vector fields $X(E)$.

By definition, a linear connection:

$$\nabla : X(E) \times X(E) \rightarrow X(E) \quad (8)$$

having the properties

$$\nabla h_T = 0, \nabla h_M = 0 \text{ and } \nabla v = 0 \quad (9)$$

is called a Γ -linear connection on the jet space of first order E .

Using the preceding definition, it immediately follows that a Γ -linear connection ∇ on the jet bundle of first order is unique determined by a set of *nine* local functions, denoted by:

$$\begin{aligned} \nabla \Gamma = & \left(\bar{G}_{\beta\gamma}^\alpha, G_{i\gamma}^k, G_{(\alpha)(j)\gamma}^{(i)(\beta)}, \bar{L}_{\beta j}^\alpha, L_{ij}^k \right. \\ & \left. L_{(\alpha)(j)k}^{(i)(\beta)}, \bar{C}_{\beta(k)}^{\alpha(\gamma)}, C_{i(k)}^{j(\gamma)}, C_{(\alpha)(j)(k)}^{(i)(\beta)(\gamma)} \right) \end{aligned} \quad (10)$$

which are defined by the relations:

$$\begin{aligned} \nabla_{\frac{\delta}{\delta t^\gamma}} \frac{\delta}{\delta t^\beta} &= \bar{G}_{\beta\gamma}^\alpha \frac{\delta}{\delta t^\alpha} \\ \nabla_{\frac{\delta}{\delta t^\gamma}} \frac{\delta}{\delta x^i} &= G_{i\gamma}^k \frac{\delta}{\delta x^k} \\ \nabla_{\frac{\delta}{\delta t^\gamma}} \frac{\partial}{\partial x_\beta^i} &= G_{(\alpha)(i)\gamma}^{(k)(\beta)} \frac{\partial}{\partial x_\alpha^k} \\ \nabla_{\frac{\delta}{\delta x^j}} \frac{\delta}{\delta t^\beta} &= \bar{L}_{\beta j}^\alpha \frac{\delta}{\delta t^\alpha} \\ \nabla_{\frac{\delta}{\delta x^j}} \frac{\delta}{\delta x^i} &= L_{ij}^k \frac{\delta}{\delta x^k} \\ \nabla_{\frac{\delta}{\delta x^j}} \frac{\partial}{\partial x_\beta^i} &= L_{(\alpha)(i)j}^{(k)(\beta)} \frac{\partial}{\partial x_\alpha^k} \\ \nabla_{\frac{\partial}{\partial x_\gamma^j}} \frac{\delta}{\delta t^\beta} &= \bar{C}_{\beta(j)}^{\alpha(\gamma)} \frac{\delta}{\delta t^\alpha} \end{aligned}$$

$$\begin{aligned}\nabla \frac{\partial}{\partial x^j_\gamma} \frac{\delta}{\delta x^i} &= C_{i(j)}^{k(\gamma)} \frac{\delta}{\delta x^k} \\ \nabla \frac{\partial}{\partial x^j_\gamma} \frac{\partial}{\partial x^i_\beta} &= C_{(\alpha)(i)(j)}^{(k)(\beta)(\gamma)} \frac{\partial}{\partial x^k_\alpha}\end{aligned}\quad (11)$$

Using the transformation laws of the elements of the adapted basis of vector fields, together with the properties of the Γ -linear connection ∇ , we deduce that the those nine adapted components $\nabla\Gamma$ transform after the rules:

$$\begin{aligned}\bar{G}_{\alpha\beta}^\delta \frac{\partial \tilde{t}^\varepsilon}{\partial t^\delta} &= \tilde{G}_{\mu\gamma}^\varepsilon \frac{\partial \tilde{t}^\mu}{\partial t^\alpha} \frac{\partial \tilde{t}^\gamma}{\partial t^\beta} + \frac{\partial^2 \tilde{t}^\varepsilon}{\partial t^\alpha \partial t^\beta} \\ G_{i\gamma}^k &= \tilde{G}_{j\beta}^m \frac{\partial x^k}{\partial \tilde{x}^m} \frac{\partial \tilde{x}^j}{\partial x^i} \frac{\partial \tilde{t}^\beta}{\partial t^\gamma} \\ G_{(\gamma)(i)\alpha}^{(k)(\beta)} &= \tilde{G}_{(\varepsilon)(j)\mu}^{(p)(\eta)} \frac{\partial x^k}{\partial \tilde{x}^p} \frac{\partial \tilde{t}^\varepsilon}{\partial t^\gamma} \frac{\partial \tilde{x}^j}{\partial x^i} \\ \frac{\partial t^\beta}{\partial \tilde{t}^\eta} \frac{\partial \tilde{t}^\mu}{\partial t^\alpha} + \delta_i^k \frac{\partial \tilde{t}^\mu}{\partial t^\alpha} \frac{\partial \tilde{t}^\varepsilon}{\partial t^\gamma} \frac{\partial^2 t^\beta}{\partial \tilde{t}^\mu \partial \tilde{t}^\varepsilon} \\ \bar{L}_{\beta j}^\gamma \frac{\partial x^j}{\partial \tilde{x}^l} &= \tilde{L}_{\mu l}^\eta \frac{\partial t^\gamma}{\partial \tilde{t}^\eta} \frac{\partial \tilde{t}^\mu}{\partial t^\beta} \\ L_{ij}^m \frac{\partial \tilde{x}^r}{\partial x^m} &= \tilde{L}_{pq}^r \frac{\partial \tilde{x}^p}{\partial x^i} \frac{\partial \tilde{x}^q}{\partial x^j} + \frac{\partial^2 \tilde{x}^r}{\partial x^i \partial x^j} \\ L_{(\gamma)(i)j}^{(k)(\beta)} &= \tilde{L}_{(v)(p)l}^{(r)(\eta)} \frac{\partial x^k}{\partial \tilde{x}^r} \frac{\partial \tilde{t}^v}{\partial t^\gamma} \frac{\partial \tilde{x}^p}{\partial x^i} \\ \frac{\partial t^\beta}{\partial \tilde{t}^\eta} \frac{\partial \tilde{x}^l}{\partial x^j} + \delta_\gamma^\beta \frac{\partial x^k}{\partial \tilde{x}^r} \frac{\partial^2 \tilde{x}^r}{\partial x^i \partial x^j} \\ \bar{C}_{\beta(i)}^{\gamma(\alpha)} &= \tilde{C}_{\varepsilon(j)}^{\mu(\delta)} \frac{\partial t^\gamma}{\partial \tilde{t}^\mu} \frac{\partial \tilde{t}^\varepsilon}{\partial t^\beta} \frac{\partial \tilde{x}^j}{\partial x^i} \frac{\partial t^\alpha}{\partial \tilde{t}^\delta} \\ C_{i(j)}^{k(\alpha)} &= \tilde{C}_{p(r)}^{s(\beta)} \frac{\partial x^k}{\partial \tilde{x}^s} \frac{\partial \tilde{x}^p}{\partial x^i} \frac{\partial \tilde{x}^r}{\partial x^j} \frac{\partial t^\alpha}{\partial \tilde{t}^\beta} \\ C_{(\gamma)(i)(j)}^{(k)(\beta)(\alpha)} &= \tilde{C}_{(\varepsilon)(p)(q)}^{(r)(\mu)(v)} \frac{\partial x^k}{\partial \tilde{x}^r} \frac{\partial \tilde{t}^\varepsilon}{\partial t^\gamma} \\ &\cdot \frac{\partial \tilde{x}^p}{\partial x^i} \frac{\partial \tilde{t}^\beta}{\partial \tilde{t}^\mu} \frac{\partial \tilde{x}^q}{\partial x^j} \frac{\partial t^\alpha}{\partial \tilde{t}^v}\end{aligned}\quad (12)$$

Finally, we point out that for a given Γ -linear connection ∇ characterized by nine adapted components $\nabla\Gamma$, we can compute the adapted components of its torsion and curvature d-tensors. In this direction, some strong results from the monograph [3] prove that the torsion d-tensor field \mathbf{T} of $\nabla\Gamma$ is determined by *twelve* effective adapted components, while the curvature d-tensor field \mathbf{R} of $\nabla\Gamma$ is determined by *eighteen* effective adapted components.

The expressions of all these effective adapted components are locally described in [3].

4. A BERWALD CONNECTION

Let us suppose now that $h_{\alpha\beta}(t)$ and $\varphi_{ij}(x)$, where $t=(t^\gamma)$ and $x=(x^k)$, are semi-Riemannian metrics on the manifolds T and M and let us consider that $H_{\alpha\beta}^\gamma(t)$ and $\Gamma_{ij}^k(x)$ are the Christoffel symbols of these metrics. Then, taking into account the transformation laws of the Christoffel symbols, by direct computations, we deduce that the pair of local functions $\Gamma_0 = (\tilde{M}_{(\alpha)\beta}^{(i)}, \tilde{N}_{(\alpha)j}^{(i)})$, where:

$$\begin{aligned}\tilde{M}_{(\alpha)\beta}^{(i)} &= -H_{\alpha\beta}^\gamma x_\gamma^i \\ \tilde{N}_{(\alpha)j}^{(i)} &= \Gamma_{jm}^\alpha x_\alpha^m\end{aligned}\quad (13)$$

represents a nonlinear connection on the 1-jet space $J^1(T, M)$.

The nonlinear connection Γ_0 is called the *canonical nonlinear connection produced by the metrics $h_{\alpha\beta}(t)$ and $\varphi_{ij}(x)$* .

Moreover, if we study the transformation laws of the following set of nine local functions, denoted by

$$B\Gamma_0 = (H_{\beta\gamma}^\alpha, 0, G_{(\alpha)(j)\gamma}^{(i)(\beta)}, 0, \Gamma_{ij}^k, L_{(\alpha)(j)k}^{(i)(\beta)}, 0, 0, 0)\quad (14)$$

where

$$G_{(\alpha)(j)\gamma}^{(i)(\beta)} = -\delta_j^i H_{\alpha\gamma}^\beta$$

$$L_{(\alpha)(j)k}^{(i)(\beta)} = \delta_{\alpha}^{\beta} \Gamma_{jk}^i \quad (15)$$

then we deduce that the adapted components $B\Gamma_0$ represent a Γ_0 -linear connection on the 1-jet space $J^1(T, M)$.

The Γ_0 -linear connection $B\Gamma_0$ is called the *Berwald connection associated to the metrics $h_{\alpha\beta}(t)$ and $\varphi_{ij}(x)$* .

Now, particularizing some general results from the work [3], we can conclude that all adapted components of the torsion d-tensor T_0 of our Berwald connection are zero, except:

$$\begin{aligned} R_{(\mu)\alpha\beta}^{(m)} &= -H_{\mu\alpha\beta}^{\gamma} x_{\gamma}^m \\ R_{(\mu)ij}^{(m)} &= r_{ijl}^m x_{\mu}^l \end{aligned} \quad (16)$$

where $H_{\mu\alpha\beta}^{\gamma}(t)$ and $r_{ijl}^m(x)$ are the classical curvature tensors of the metrics $h_{\alpha\beta}(t)$ and $\varphi_{ij}(x)$.

Also, all adapted components of the curvature d-tensor R_0 of our Berwald connection are zero, except:

$$R_{\alpha\beta\gamma}^{\delta} = H_{\alpha\beta\gamma}^{\delta}, \quad R_{ijk}^l = r_{ijk}^l \quad (17)$$

5. CONCLUSIONS

As final remarks, let us note that, in the particular case $(T, h) = (R, \delta)$, the canonical nonlinear connection Γ_0 naturally generalizes

the canonical nonlinear connection produced by the spray $2G^i = \Gamma_{jk}^i y^j y^k$, while our Berwald connection is the natural generalization of the classical linear Berwald connection from the theory of Lagrange spaces [2], which is produced by the nonlinear connection $N_j^i = \Gamma_{jk}^i y^k$.

In conclusion, the Berwald Γ_0 -linear connection $B\Gamma_0$ is a natural example of Γ -linear connection, which ensures us that our geometrical theory upon the Γ -linear connections on 1-jet spaces is a fertile and good one.

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